

# PVDIS Physics Case

# Main Bullets

- BSM Physics via  $C_2$ 's
- CSV in deuterium
- CSV in  $^{48}\text{Ca}$ ?
- d/u in the proton

# Perspective

- Important new results in the field have been released: LHC, Qweak.
- Physics motivation is more important for Science Review than PCDR.
- We plan to produce updated document about the SoLID PVDIS physics case.

# Standard Model Test: Couplings

This has  
not changed

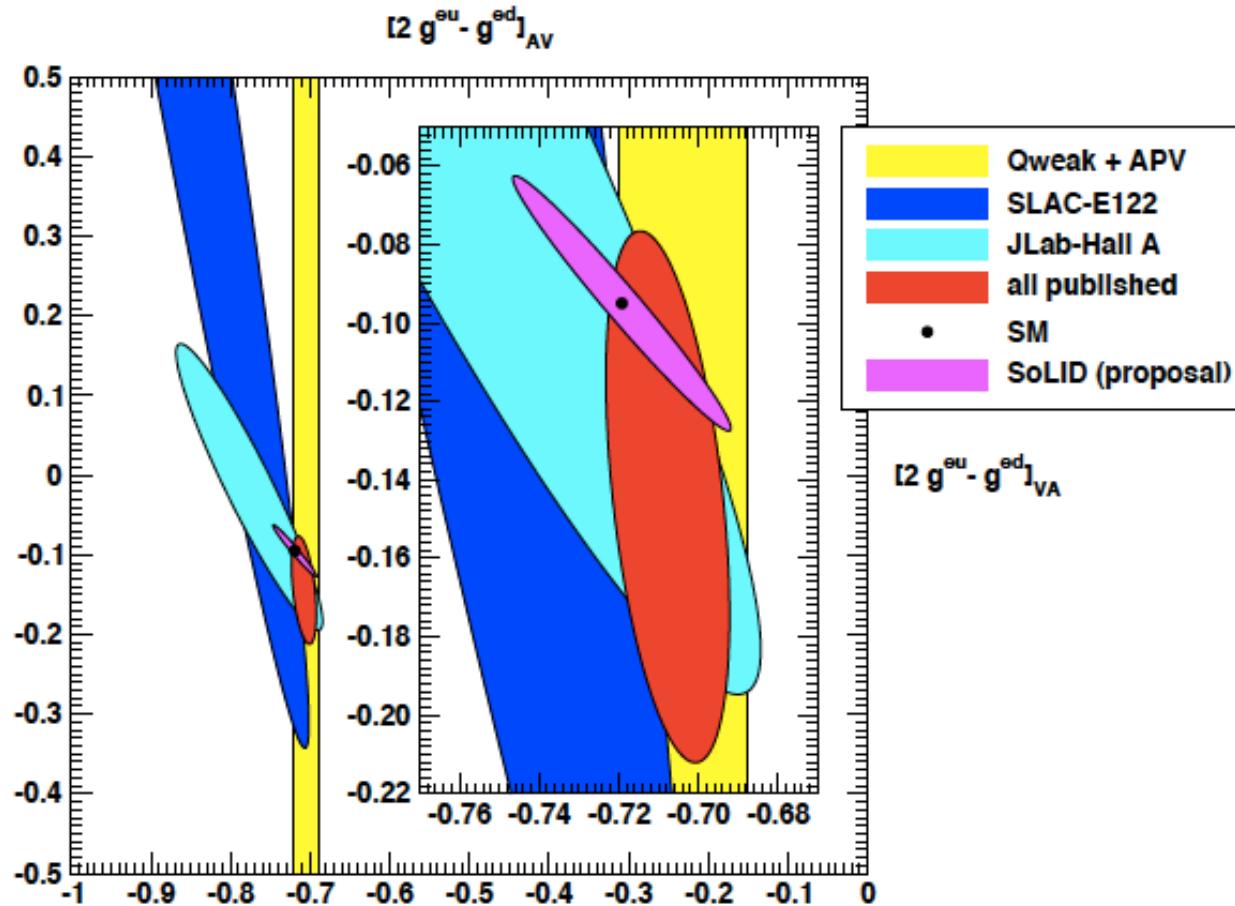


Figure 7: Results form the JLab PVDIS collaboration together with the projected results from the SoLID PVDIS experiment. SoLID Collaboration: PVDIS

# LHC $\text{pp} \rightarrow e^+e^-$ and SoLID PVDIS

- LHC has confirmed no new physics:  $Z'$ , etc.
- Both measure the same contact interactions.
- For the LHC,  $Q^2$  is  $10^5\text{-}10^6$  times larger.
- LHC systematics are large:  $>10\%$ .
- Much LHC sensitivity comes where square of contact interactions are significant.
- LHC sets bounds on sum of squares of all couplings.

# Contact Interactions: Compositeness

CI Lagrangian:

$$\mathcal{L}_{eff} = \frac{g^2}{\Lambda^2} \sum_{i,j=L,R} \eta_{ij}^f \bar{e}_i \gamma_\mu e_i \bar{f}_j \gamma^\mu f_j$$

$e^+e^-$  at LHC:

$$\frac{d\sigma}{d\Omega} (q_L \bar{q}_R \rightarrow e_L^- e_R^+) = \frac{\alpha^2}{4s} (1 + \cos\theta)^2 |Q_q - r L_q L_e - \zeta \eta_{LL}^{qe}|^2$$

$$r = \frac{1}{4 \sin^2 \theta_W \cos^2 \theta_W} \sim 1.4$$

$$\zeta = \frac{M^2}{\alpha \Lambda^2} \sim 1 \text{ for LHC}$$

	$e^-$	$u$	$d$
$Q$	-1	0.67	-0.33
$L$	-0.53	0.69	-0.84
$R$	0.47	-0.31	0.16

At large  $M$ ,  $\sigma_T \sim \zeta^2 (\eta_{LL}^2 + \eta_{LR}^2 + \eta_{RR}^2 + \eta_{RL}^2)$  Sensitive to  $C_2$ 's

# Comment From ATLAS

The relative impact of the interference and pure CI terms depends on both the dilepton mass and  $\Lambda$ . For example, the magnitude of the interference term for dilepton masses above 600 GeV is about twice as large as that of the pure CI term at  $\Lambda = 14$  TeV; the interference becomes increasingly dominant for higher values of  $\Lambda$ .

# ATLAS 8 TeV

Process	$m_{ee}$ [GeV]					
	400–550	550–800	800–1200	1200–1800	1800–3000	3000–4500
Total SM	$1260 \pm 100$	$404 \pm 35$	$82 \pm 9$	$10.8 \pm 1.6$	$0.91 \pm 0.21$	$0.014 \pm 0.005$
Data	1262	388	84	7	0	0
SM+CI ( $\Lambda_{LL}^- = 14$ TeV)	$1310 \pm 110$	$440 \pm 40$	$108 \pm 10$	$20.9 \pm 1.9$	$4.2 \pm 0.4$	$0.141 \pm 0.028$
SM+CI ( $\Lambda_{LL}^- = 20$ TeV)	$1290 \pm 110$	$430 \pm 40$	$90 \pm 10$	$14.4 \pm 1.7$	$2.01 \pm 0.25$	$0.045 \pm 0.012$
SM+CI ( $\Lambda_{LR}^- = 14$ TeV)	$1340 \pm 110$	$460 \pm 40$	$118 \pm 10$	$26.3 \pm 2.1$	$6.0 \pm 0.5$	$0.28 \pm 0.05$
SM+CI ( $\Lambda_{LR}^- = 20$ TeV)	$1290 \pm 110$	$420 \pm 40$	$98 \pm 10$	$15.7 \pm 1.7$	$2.58 \pm 0.28$	$0.078 \pm 0.018$

Contact Interactions appear above  
systematics only at large  $M_{ee}$

# ATLAS 13 TeV

$m_{ee}$ [GeV]	80–120	120–250	250–400	400–500	500–700
Drell–Yan	$11\,800\,000 \pm 700\,000$	$216\,000 \pm 11\,000$	$17\,230 \pm 1000$	$2640 \pm 180$	$1620 \pm 120$
Top quarks	$28\,600 \pm 1800$	$44\,600 \pm 2900$	$8300 \pm 600$	$1130 \pm 80$	$560 \pm 40$
Dibosons	$31\,400 \pm 3300$	$7000 \pm 700$	$1300 \pm 140$	$228 \pm 25$	$146 \pm 16$
Multi-jet & $W+jets$	$11\,000 + 9000$	$5600 \pm 2000$	$780 \pm 80$	$151 \pm 21$	$113 \pm 17$
Total SM	$11\,900\,000 \pm 700\,000$	$273\,000 \pm 12\,000$	$27\,600 \pm 1100$	$4150 \pm 200$	$2440 \pm 130$
Data	$12\,415\,434$	$275\,711$	$27\,538$	$4140$	$2390$
$Z'_\chi$ (4 TeV)	$0.00635 \pm 0.00021$	$0.0390 \pm 0.0015$	$0.0564 \pm 0.0025$	$0.0334 \pm 0.0027$	$0.064 \pm 0.004$
$Z'_\chi$ (5 TeV)	$0.00305 \pm 0.00012$	$0.0165 \pm 0.0006$	$0.0225 \pm 0.0010$	$0.0139 \pm 0.0007$	$0.0275 \pm 0.0015$
$m_{ee}$ [GeV]	700–900	900–1200	1200–1800	1800–3000	3000–6000
Drell–Yan	$421 \pm 34$	$176 \pm 17$	$62 \pm 7$	$8.7 \pm 1.3$	$0.34 \pm 0.07$
Top quarks	$94 \pm 8$	$27.9 \pm 2.8$	$5.1 \pm 0.7$	$< 0.001$	$< 0.001$
Dibosons	$39 \pm 4$	$16.9 \pm 2.1$	$5.8 \pm 0.8$	$0.74 \pm 0.11$	$0.028 \pm 0.004$
Multi-jet & $W+jets$	$39 \pm 6$	$16.1 \pm 2.0$	$7.9 \pm 2.3$	$1.6 \pm 1.2$	$0.08 \pm 0.27$
Total SM	$590 \pm 40$	$237 \pm 17$	$81 \pm 7$	$11.0 \pm 1.8$	$0.45 \pm 0.28$
Data	$589$	$209$	$61$	$10$	$0$
$Z'_\chi$ (4 TeV)	$0.0585 \pm 0.0035$	$0.074 \pm 0.005$	$0.121 \pm 0.011$	$0.172 \pm 0.017$	$2.57 \pm 0.27$
$Z'_\chi$ (5 TeV)	$0.0218 \pm 0.0013$	$0.0295 \pm 0.0021$	$0.040 \pm 0.004$	$0.040 \pm 0.004$	$0.280 \pm 0.030$

# Obsolete Plot?

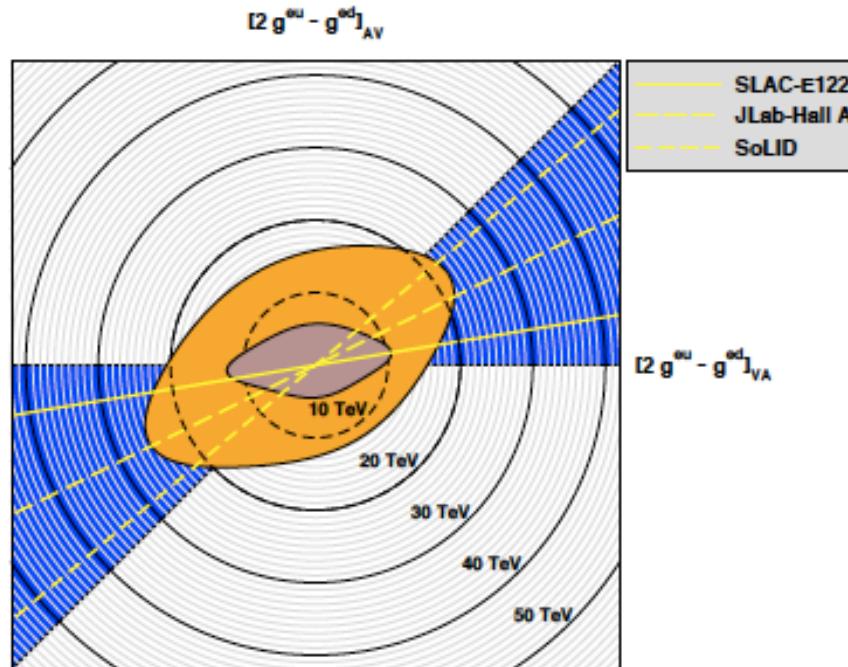
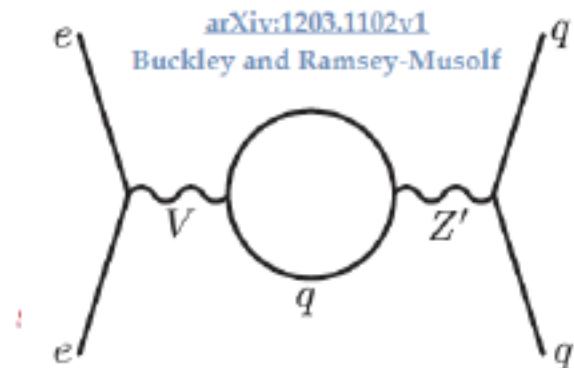


Figure 10: Polar plot for limits on composite models. The gray ellipse includes the published results from the PVDIS and Qweak collaborations. The orange ellipse gives the projected limits with the full Qweak statistics and the SoLID data.

# Physics Unaffected by the LHC

## *Leptophobic Z'*

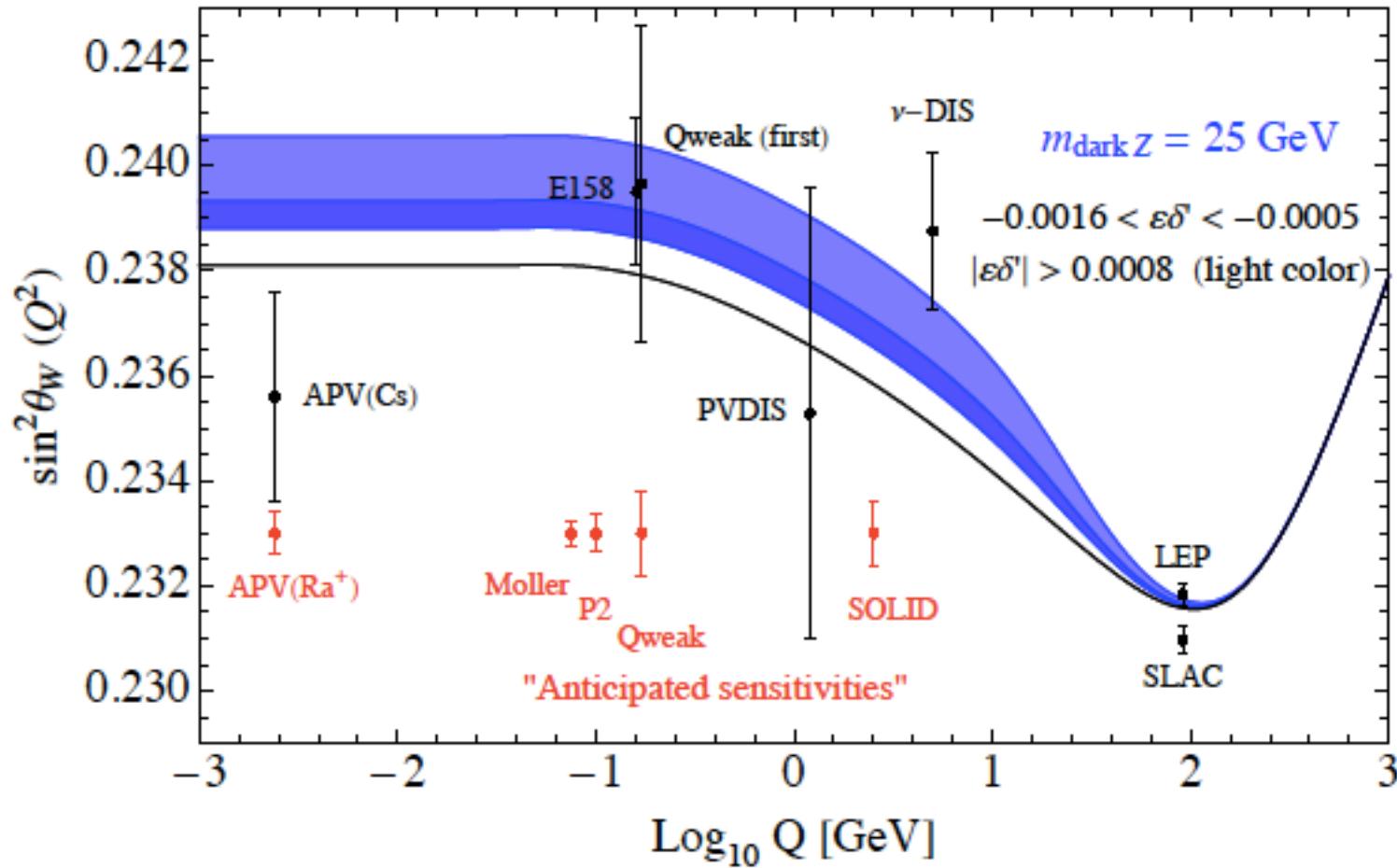


SOLID can improve sensitivity:  
100-200 GeV range

ATLAS rules out  
Leptophobic  $Z'$   
between  
1.1 and 1.5 TeV

Figure 9: Diagram of a leptophobic  $Z'$  that can contribute to the  $C_{2l}$  and few other observables.

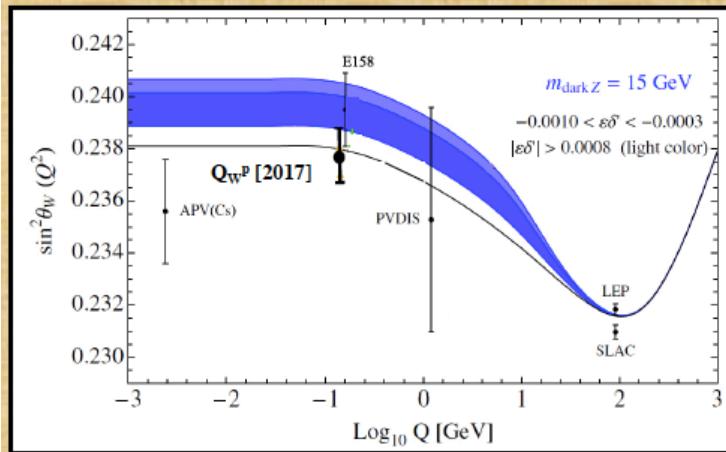
# New Idea: Dark Light



# New Results from Qweak

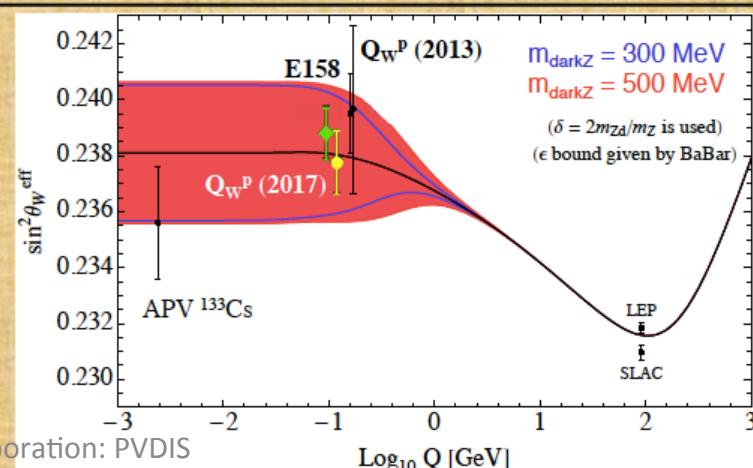
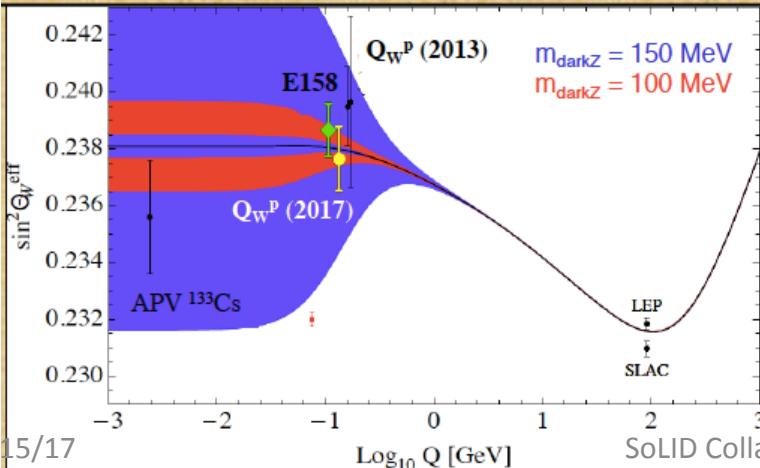
## Implications for “Dark Parity Violation”

“Dark photon” – possible portal for new force to communicate with SM?



(Davoudiasl, Lee, Marciano, arXiv 1402.3620)

- Astrophysical motivation: observed excess in positron data.
- Introduces new source of low energy parity violation through mass mixing between  $Z$  and  $Z_d$  with observable consequences.
- Complementary to direct searches for heavy dark photons.



# $d/u$ : Old

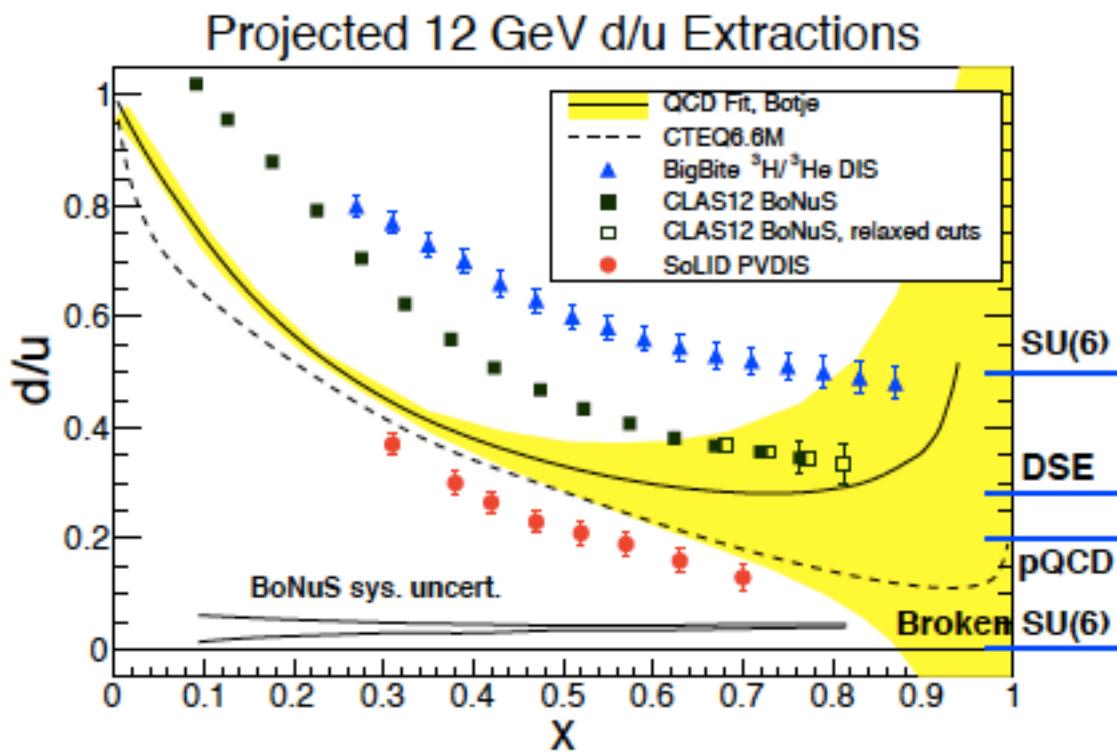
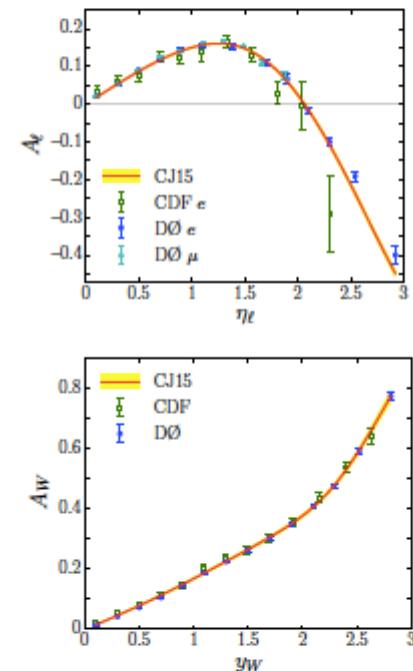
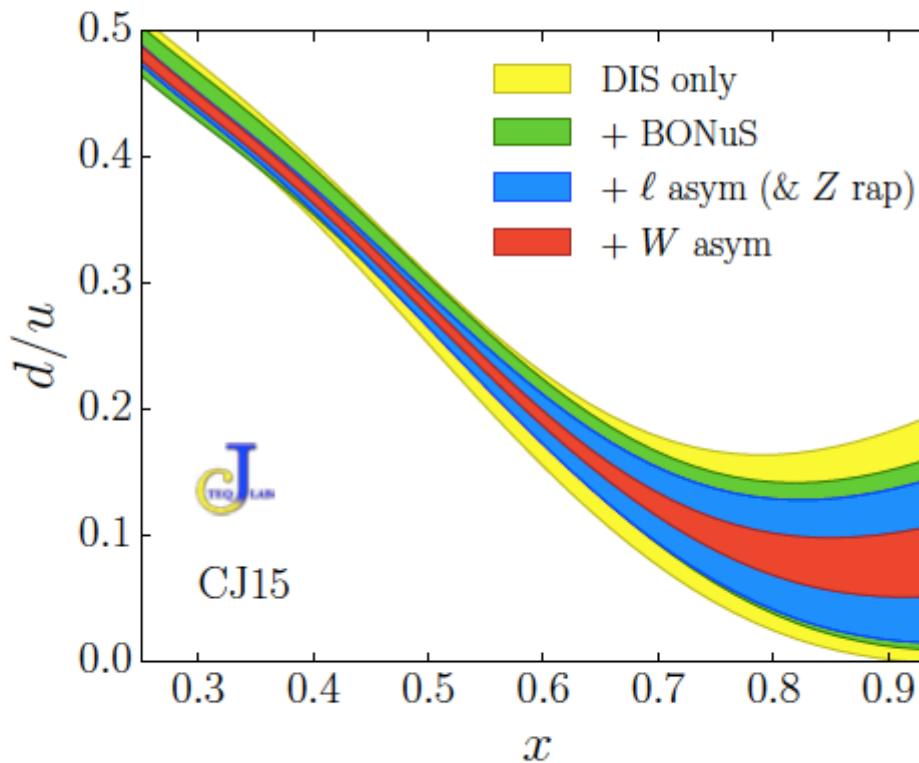


Figure 13: Anticipated precision for  $d/u$  measurement with Solid as well as other proposed experiments.  
10/15/17  
SoLID Collaboration: PVDIS

## ■ At large $x$ proton asymmetry sensitive to $d/u$ ratio

$$a_1 = \frac{12C_{1u} - 6C_{1d} d/u}{4 + d/u}$$

$$a_3 = \frac{12C_{2u} - 6C_{2d} d/u}{4 + d/u}$$



- error reduced by recent  $W$  asymmetry data
  - but rely on reconstruction from  $W \rightarrow \ell \nu$  decays

# 1 Measurement of $d/u$ for the Proton

## 1.1 Main Bullets

1. Interplay and complementary of the different experiments: There are four complementary precision experiments:

- (a) PVDIS is the most direct measurement of  $d/u$
- (b) BONUS: some nuclear corrections.
- (c) Marathon: short-range correlations.
- (d) Fermilab collider: less direct,  $Q^2 \sim M_W^2$

PVDIS provides low  $Q^2$  measure of  $d/u$  free of any nuclear corrections.  
It is the most direct method.

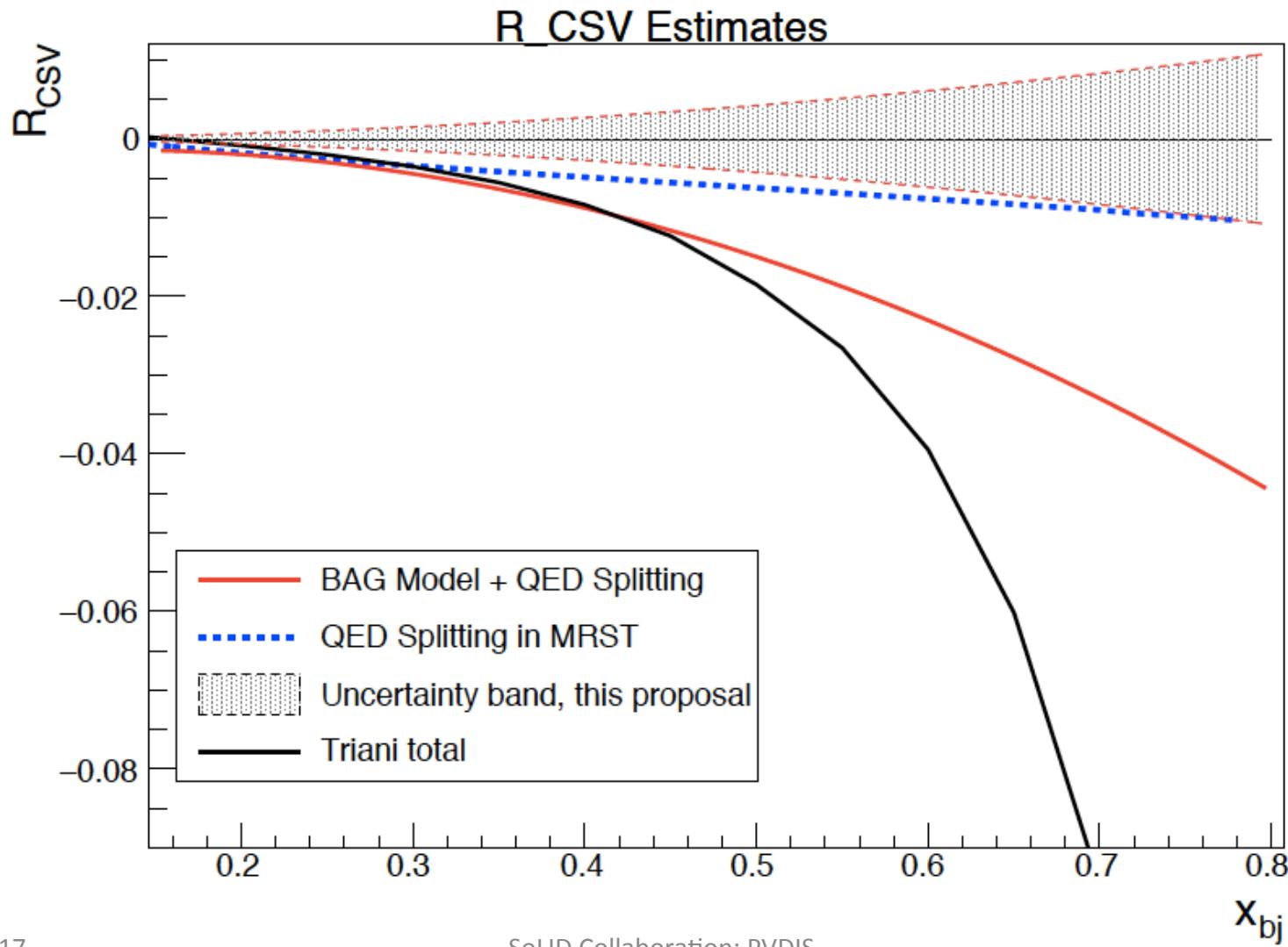
## 2. Possible Plots

- (a) Plot showing impact of projected PVDIS data on global fits for  $d/u$ .
- (b) Show experiments on the plot in the CJ15 paper showing nuclear correction for deuteron.
- (c) Plot showing Fermilab collider data on the  $d/u$  plot?

## 3. Physics impact of $d/u$ :

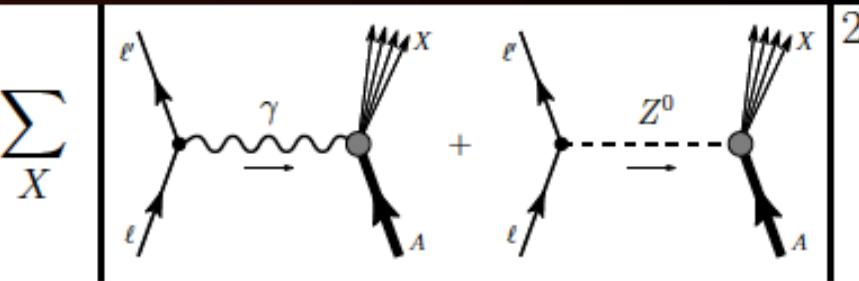
- (a) Quark models as  $x \rightarrow 1$

# CSV: Found New Model



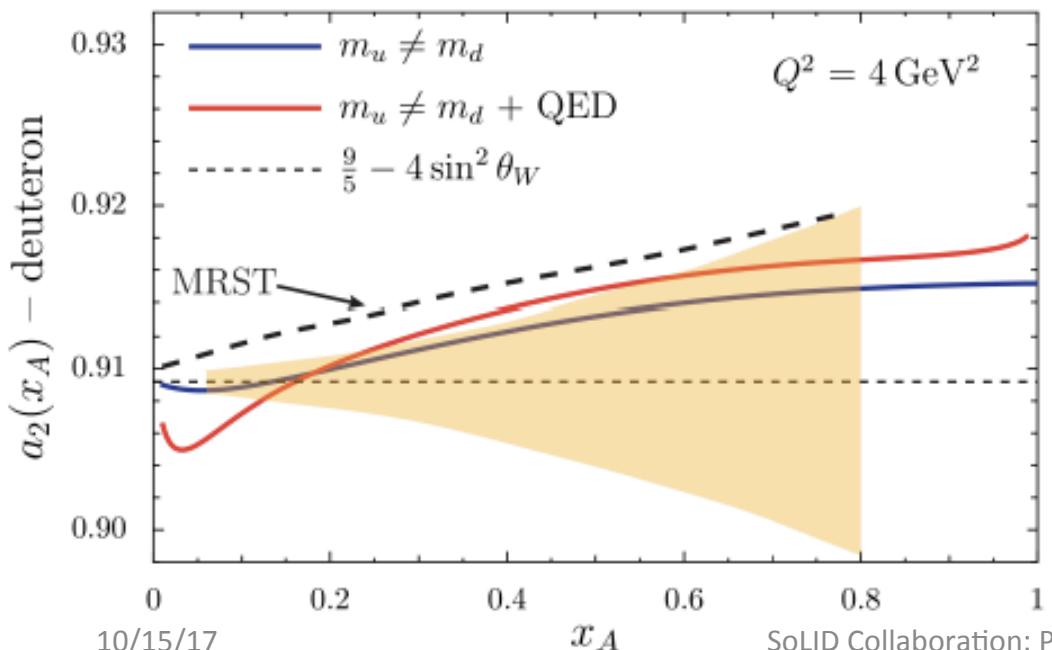
# Measuring CSB using PVDIS

- Consider the deuteron where nuclear effects are well known
- Construct the PVDIS asymmetry



$$a_2(x) \stackrel{N \sim Z}{=} \frac{9}{5} - 4 \sin^2 \theta_W - \frac{12}{25} \frac{u_D^+(x) - d_D^+(x)}{u_D^+(x) + d_D^+(x)}; \quad \Delta a_2 = -\frac{6}{25} \frac{\delta u^+ - \delta d^+}{\delta u^+ + \delta d^+}$$

- PVDIS on  $N = Z$  target – excellent way to access CSB in nucleon PDFs



- Our result likely too small to be measured in proposed JLab 12 GeV experiment
- QED splitting does increase CSB effects but not significantly

# New CSB Results from Quark Masses

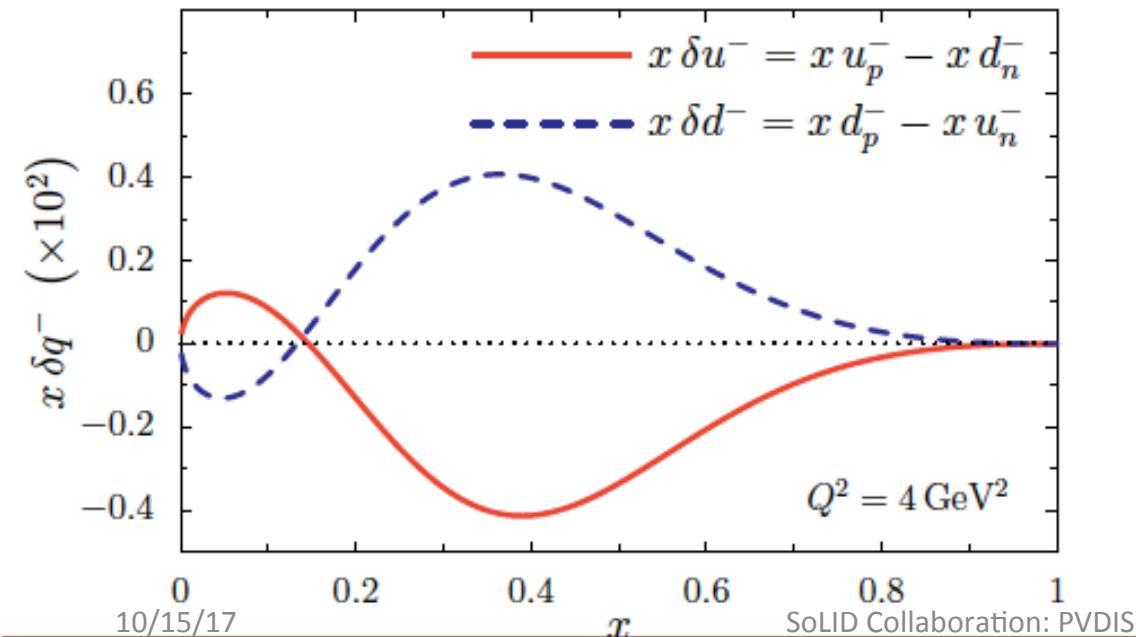
- Our dressed quark mass is  $M_0 = 400 \text{ MeV}$  in the charge symmetry limit
  - define:  $M_u = M_0 - \delta_M$  &  $M_d = M_0 + \delta_M$ ;  $\delta_M \mid M_n - M_p = 3 \text{ MeV}$

$$M_u = 397 \text{ MeV} \quad \& \quad M_d = 403 \text{ MeV} \implies \frac{M_d - M_u}{M_d} \simeq 1.5\%; \quad \frac{m_u}{m_d} = 0.58$$

- For the CSB PDFs we find, at  $Q^2 = 4 \text{ GeV}^2$ :

$$\langle x \delta u^- \rangle = -0.0019 \quad \langle x \delta d^- \rangle = 0.0019$$

- therefore a small amount of momentum is shifted from the  $u$  to the  $d$  quarks



- Recent lattice QCD analysis at  $Q^2 = 4 \text{ GeV}^2$  find:

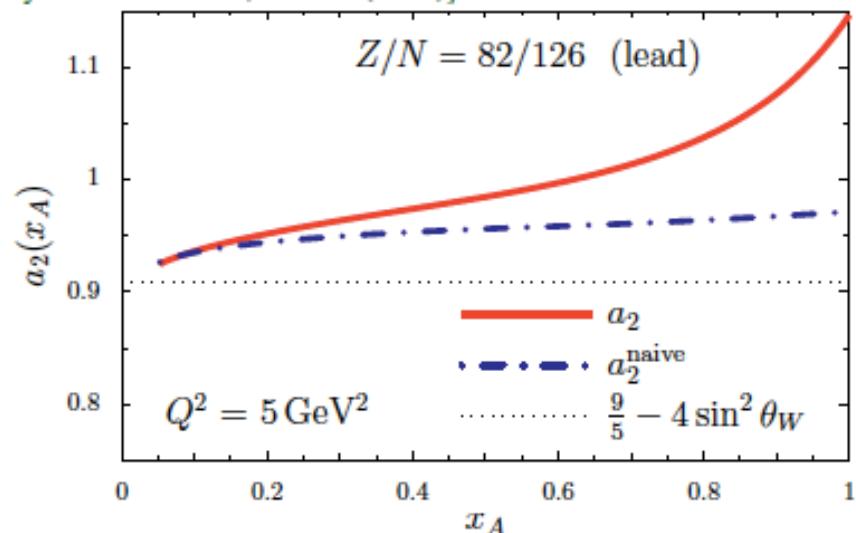
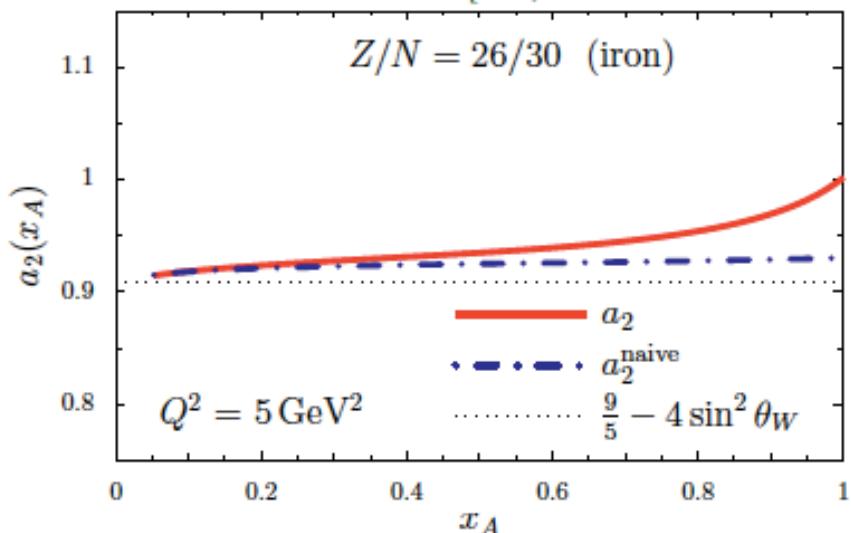
[R. Horsley *et al.*, PRD 83, 051501 (2011)]

$$\begin{aligned}\langle x \delta u \rangle &= -0.0023(6) \\ \langle x \delta d \rangle &= 0.0020(3)\end{aligned}$$

- Therefore our results are consistent with Lattice

# Isovector Effects in Nuclei

[ICC, W. Bentz and A. W. Thomas, Phys. Rev. Lett. 109, 182301 (2012)]



- PVDIS –  $\gamma Z$  interference:

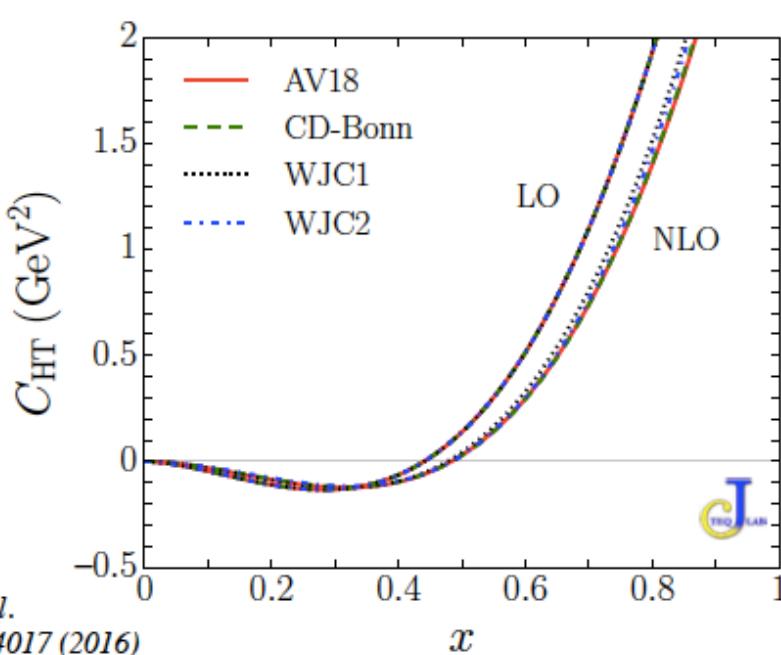
$$a_2(x) = -2 g_A^e \frac{F_2^{\gamma Z}(x)}{F_2^\gamma(x)} \stackrel{N \sim Z}{\simeq} \frac{9}{5} - 4 \sin^2 \theta_W - \frac{12}{25} \frac{u_A^+(x) - d_A^+(x)}{u_A^+(x) + d_A^+(x)}$$

- Deviation from naive expectation: momentum shifted *from u to d quarks*
- $F_2^{\gamma Z}(x)$  has markedly different flavour dependence compared with  $F_2^\gamma(x)$ 
  - a measurement of both enables an extraction of  $u(x)$  and  $d(x)$  separately
- **Proposal to measure  $a_2(x)$  of  $^{48}\text{Ca}$  was deferred twice ...**

# Dynamical higher twists

- In global QCD analyses parametrize structure function as

$$F_2(x, Q^2) = F_2^{(\text{LT})}(x, Q^2) \left( 1 + \frac{C_{\text{HT}}(x)}{Q^2} \right)$$



$$C_{\text{HT}} = h_0 x^{h_1} (1 + h_2 x)$$

→ cannot accommodate  $Q^2$  dependence of data without power corrections

- crucial for extracting correct PDFs at high  $x$
- can do the same with PVDIS data, especially with sufficient  $Q^2$  coverage – isolate isospin dependence of higher twists vs  $x$ !

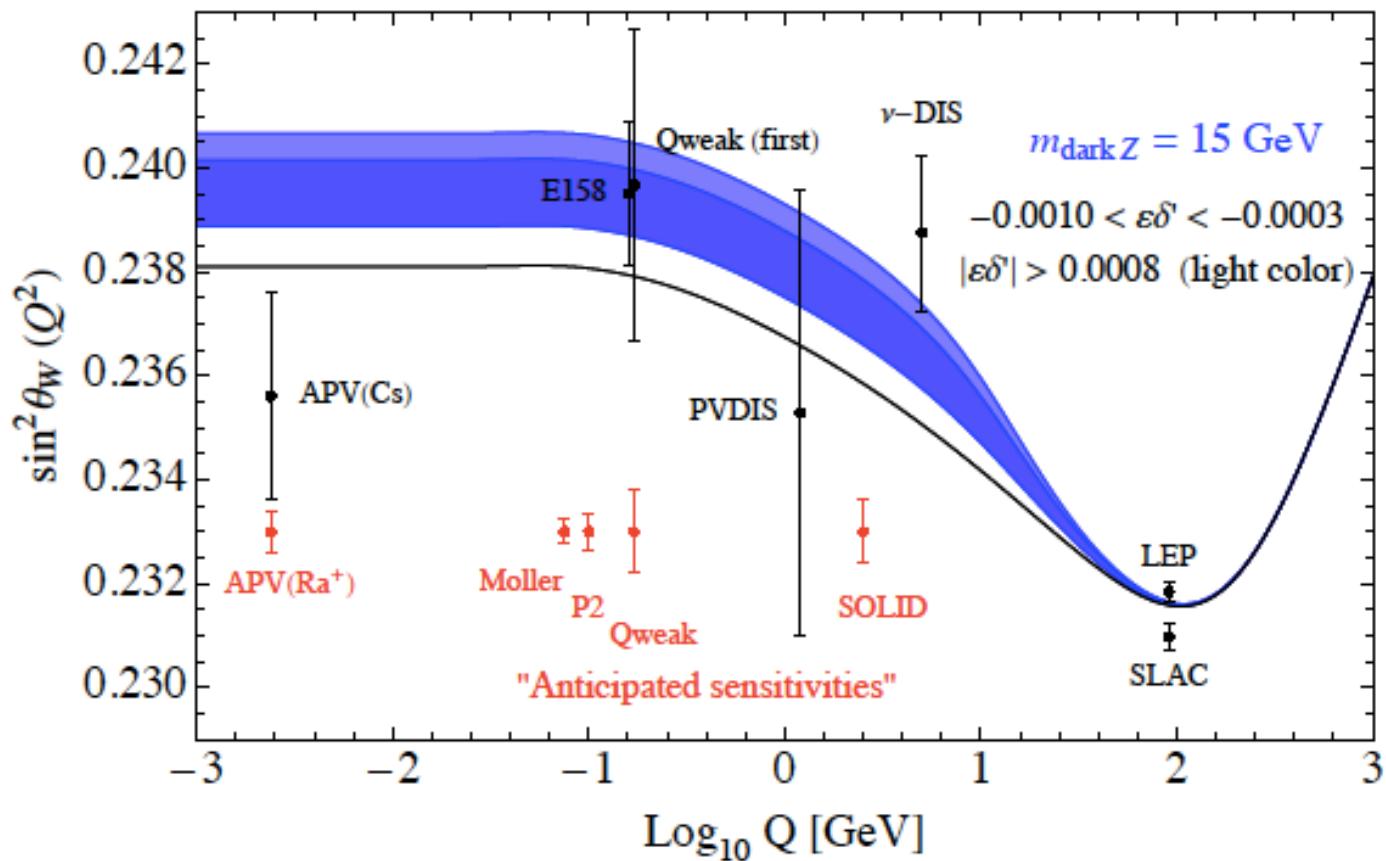


FIG. 1: Running of  $\sin^2 \theta_W(Q^2)$  with  $Q^2$  due to a  $Z_d$  with  $m_{Z_d} = 15$  GeV for a 1-sigma fit to  $\varepsilon\delta'$  from Eq. (12), shown here as the blue band. The lighter part of the band corresponds to parameters that are in some tension with precision constraints indicated by Eq. (15) (see text for details). The black curve represents the running of  $\sin^2 \theta_W(Q^2)$  in the SM [14–17]. The sizes of the red error bars only represent the anticipated sensitivities of various experiments. This figure is taken from Ref. [1].